

Coupling fields and underlying space curvature: An augmented Lagrangian approach

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We demonstrate a systematic implementation of coupling between a scalar field and the geometry of the space which carries the field. This naturally gives rise to a feedback mechanism between the field and the geometry. We develop a systematic model for the feedback in a general form, inspired by a specific implementation in the context of molecular dynamics (the so-called Rahman-Parrinello molecular dynamics, or RP-MD). We use a generalized Lagrangian that allows for the coupling of the space's metric tensor to the scalar field, and add terms motivated by RP-MD. We present two implementations of the scheme: one in which the metric is only time-dependent (which gives rise to an ordinary differential equation for its temporal evolution), and the other with spatiotemporal dependence (wherein the metric's evolution is governed by a partial differential equation). Numerical results are reported for the $(1+1)$ -dimensional model with a nonlinearity of the sine-Gordon type.

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Recently, much attention has been focused on soft-condensed-matter objects, such as vesicles, microtubules, and membranes [1–4]. Many nanoscale physical systems, including nanotubes and electronic and photonic waveguide structures [5,6], have nontrivial geometry and are influenced by substrate effects. These classes of systems, many of which are inherently nonlinear, raise the question of the interplay between nonlinearity and a substrate with variable curvature. Of particular interest is a possibility of developing curvature in the substrate due to forces generated by the nonlinear field. The resulting curvature can in turn affect the field.

There is an increasing body of literature dealing with the interplay of nonlinearity and a curved substrate. Usually, however, the substrate geometry is assumed to be *fixed*, see, e.g., Ref. [7]. Nevertheless, for many applications it is relevant to introduce models that admit a flexible substrate, which is affected by the field(s) that it carries, as well as feeding back into the field dynamics. The equations should include both the field dynamics proper and the feedback coupling to the substrate. Equations for the evolution of the substrate should in turn be affected by the evolution of the field. A prototypical physical example of this type is Euler buckling [8], where the evolution of a stress profile causes the underlying surface to buckle (and hence locally modify its curvature).

In a discrete setting, a model of this type has recently been presented in Ref. [9]. Some studies have also been performed in a special case of the continuum limit of classical spin systems (such as the Heisenberg chain) coupled to the curvature; geometric frustration was found to arise in such settings [10].

About 20 years ago, a problem similar to the theme of our study was examined in the context of molecular dynamics modelling of structural transitions in crystals. In particular, in a series of papers [11], Rahman and Parrinello introduced a new idea for studying such transitions by means of an augmented Lagrangian that would account for the degrees of freedom of the “box” (the cell), to which MD is confined. In

studying the time evolution of the box dynamics (naturally obtained through the Euler-Lagrange equations of the augmented Lagrangian for the box degrees of freedom), they were able to identify structural transitions (under external shear) from square to hexagonal patterns, fcc to bcc, etc. In the same spirit, a Lagrangian containing the curvature of a slab as an additional degree of freedom was introduced for studying strained surfaces in Ref. [12]. We will hereafter refer to this as the RP-MD (Rahman-Parrinello molecular dynamics) method. The Lagrangian for the particles and the box in this case reads

$$L = \frac{1}{2} \sum_i m_i \dot{\mathbf{s}}_i G \dot{\mathbf{s}}_i - \sum_{i,j>i} V(r_{ij}) + \frac{1}{2} W \text{Tr}(\dot{f}^T \dot{f}), \quad (1)$$

where m_i is the mass of the i th particle, $\dot{\mathbf{s}}_i$ is its vectorial velocity, the spatial part G of the spatiotemporal metric tensor may be represented in terms of another matrix f as $G = f^T f$ (G is positive definite). T and Tr denote the transposition and trace respectively, r_{ij} and V are distances between the particles and the potential of interaction between them, and W is an effective mass of the box.

Our purpose in this work is to extend the RP-MD methodology to the case of a continuum scalar field, coupled to either a spatially averaged geometric characteristic (“average curvature”), which will give rise to an ODE (ordinary differential equation), or to a spatiotemporal curvature field that will generate a PDE (partial differential equation). The continuum field may represent, e.g., a chemical concentration propagating over a membrane, or in a salt solution, or an envelope wave of the electric field in nanosystems. The spatiotemporal metric is assumed to have the simple form

$$g = \begin{pmatrix} -1 & 0 \\ 0 & G \end{pmatrix}. \quad (2)$$

We will first consider the general case, where G is a $d \times d$ matrix, d being the space dimension.

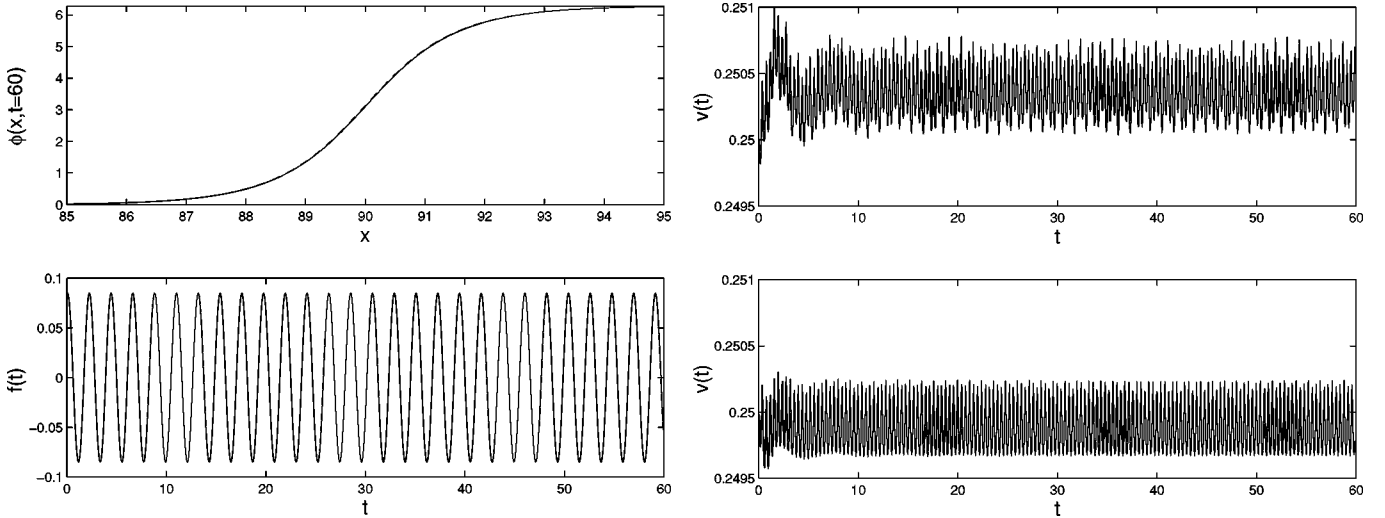


FIG. 1. The top left panel shows the kink's spatial profile at $t=60$ in model A. The initial kink profile (centered around the new center position) is shown by the dashed line, and is practically indistinguishable from the solution at $t=60$. The bottom left panel shows smooth oscillations of $f(t)$. The top and bottom right panels show, respectively, the kink's velocity $v(t)$ vs t , and the same quantity but for $f=0$.

One can define a field-type generalization of the RP-MD model, with a scalar field ϕ , as follows:

$$L = \int d^d x \left[-g^{ij} \frac{\partial \phi}{\partial x^i} \frac{\partial \phi}{\partial x^j} - V(\phi) \right] + \frac{1}{2} W \text{Tr}(f^T f), \quad (3)$$

where $V(\phi)$ is a potential governing the nonlinear evolution of the field ϕ and $f=f(t)$. If f (and hence G) is a function of both spatial coordinates and time, an elastic-energy term [13] should be added to the Lagrangian (3), so that it becomes

$$L = \int d^d x \left[-g^{ij} \frac{\partial \phi}{\partial x^i} \frac{\partial \phi}{\partial x^j} - V(\phi) \right] + \int d^d x \left\{ \frac{1}{2} W \left[\text{Tr}(f^T f) - \frac{1}{2} \text{Tr} \left(\frac{\partial f^T}{\partial x_i} \frac{\partial f}{\partial x^i} \right) \right] \right\}. \quad (4)$$

In this Brief Report, we propose two models that include the feedback to the curvature, and, simultaneously, admit as particular solutions those with a flat metric. In particular, we choose $G=1+f^2$ (so that the metric is positive definite and for $f=0$ has the original Minkowskian metric as a special case). Notice a deviation in our choice from the RP case of $G=f^2$ [14]. Equations (3) and (4) are general; we hereafter study the $d=1$ case.

Assuming initially that $f=f(t)$ (e.g., including only an effect of the ‘‘mean curvature’’ on the scalar-field dynamics), the Lagrangian (3) becomes

$$L = \frac{1}{2} W \dot{f}^2 + \int dx \left[\frac{1}{2} \left(\frac{\partial \phi}{\partial t} \right)^2 - \frac{1+f^2}{2} \left(\frac{\partial \phi}{\partial x} \right)^2 - V(\phi) \right]. \quad (5)$$

The resulting equations of motion (model A) are

$$\phi_{tt} = (1+f^2) \phi_{xx} - (\partial V / \partial \phi), \quad (6)$$

$$W f_{tt} = -f \int \phi_x^2 dx, \quad (7)$$

where the subscripts stand for the partial derivatives.

Notice that the function f is directly related to the scalar curvature of the 1D space. In particular, the Ricci scalar, which is $R=2R_{1212}/\det(g)$ [16] in the general case, in the 1D case is $R=-2\ddot{f}_{tt}/\ddot{f}$, where $\ddot{f}=\sqrt{1+f^2}$.

On the other hand, for a metric with both spatial and

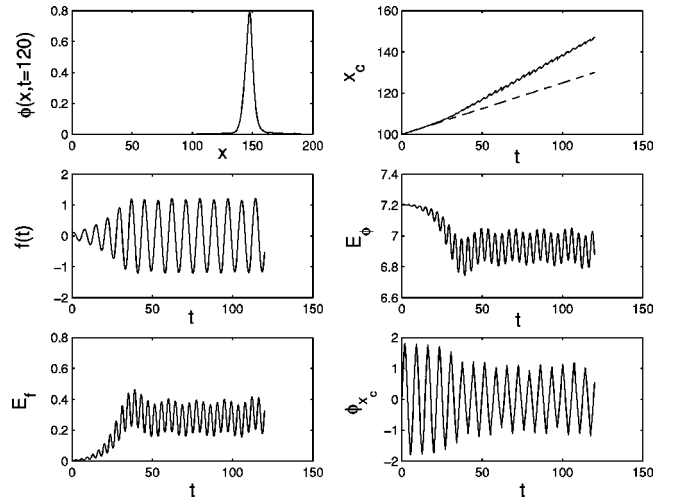


FIG. 2. The top left and right panels show, respectively, the breather in model A at the end of the simulation period, $t=120$, and its position vs time (solid), as compared to that which it would have moving at the initial velocity, $v=0.25$ (dashed). The middle left panel shows the time evolution of $f(t)$, while the middle right and bottom left panels present exchange between the curvature-mode's energy $E_f=f_t^2/2+(f^2/4)\int \phi_x^2 dx$ and the rest of the energy $E_\phi=E-E_f$. Finally, the bottom right panel shows the time evolution of the field ϕ at the center of the breather.

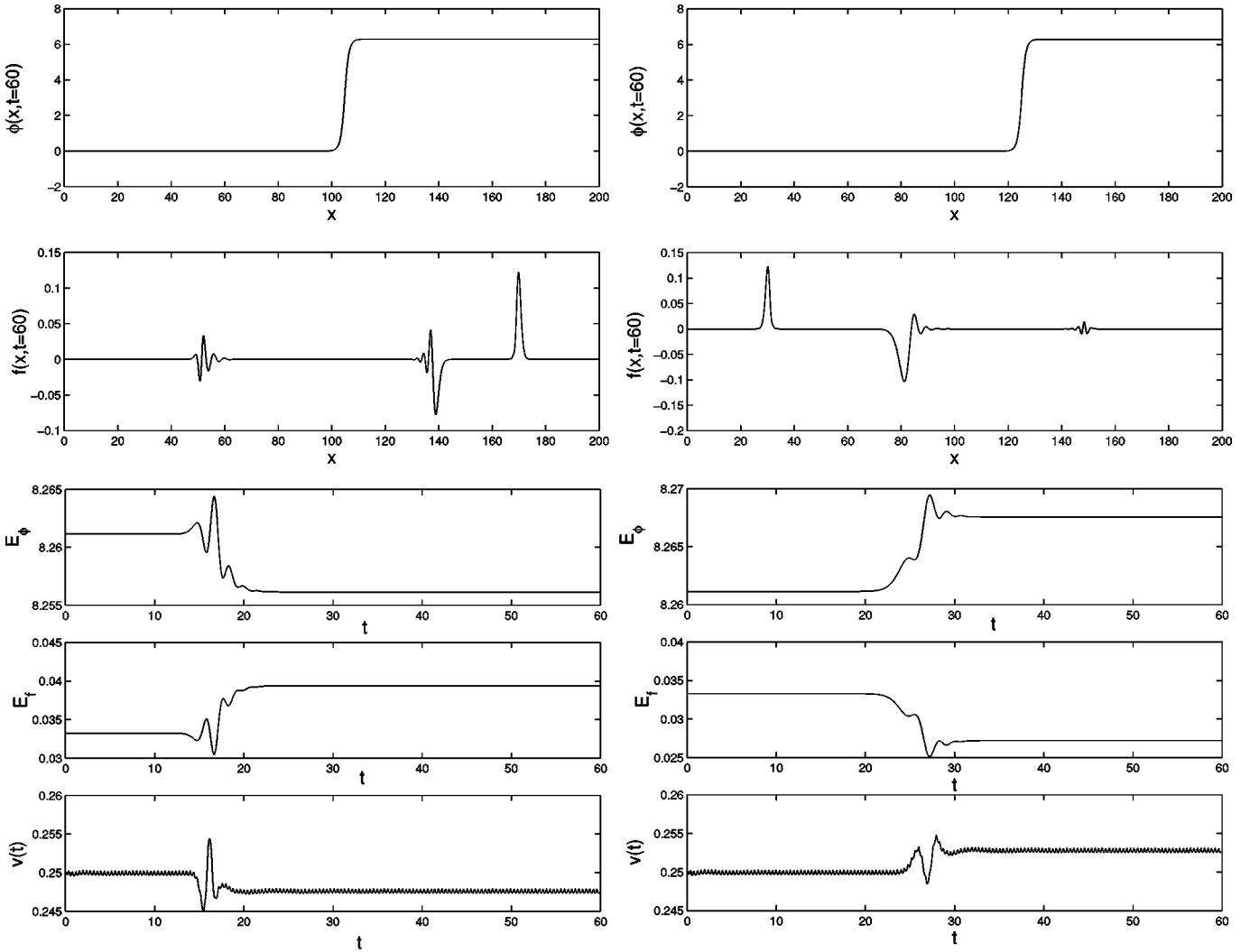


FIG. 3. The left part of the figure shows spatial profiles of the fields ϕ and f at $t=60$ in model B . It also shows the time evolution of the energies $E_\phi = E - E_f$ and $E_f = \int (1/2)[f_t^2 + f_x^2 + (f\phi_x)^2/2]dx$, as well as the kink's velocity as a function of time. The kink's center is initially at $x=90$, while the center of the curvature pulse is at $x=110$. The right subplots show counterparts of those in the left part, but for the case of the kink and curvature pulse initially centered at $x=110$ and $x=90$, respectively.

temporal dependence [e.g., for $f=f(x,t)$], one arrives at the following Lagrangian:

$$L = \int dx \left[\frac{1}{2} \left(\frac{\partial \phi}{\partial t} \right)^2 - \frac{1+f^2}{2} \left(\frac{\partial \phi}{\partial x} \right)^2 - V(\phi) \right] + \int dx \left[\frac{W}{2} \left(\frac{\partial f}{\partial t} \right)^2 - \frac{W}{2} \left(\frac{\partial f}{\partial x} \right)^2 \right]. \quad (8)$$

The ensuing equations (referred to as model B) are

$$\phi_{tt} = [(1+f^2)\phi_x]_x - (\partial V/\partial \phi), \quad (9)$$

$$Wf_{tt} = Wf_{xx} - f\phi_x^2. \quad (10)$$

As a particular application of models A and B , we examine the physically ubiquitous sine-Gordon (sG) potential, $V(\phi) = 1 - \cos(\phi)$ [15]. It is clear that Eqs. (6) and (7) and (9) and (10) have particular solutions with $f=0$, for which

the second equation in each pair is satisfied trivially, while the first equation reduces to the sG equation. Basic solitary-wave solutions of the sG equation are the topological soliton (kink),

$$\phi_k(x,t) = 4 \tan^{-1} \{ \exp[\gamma(x-x_0-vt)] \}, \quad (11)$$

where v is its velocity, $\gamma = (1-v^2)^{-1/2}$ is the Lorentz factor, and x_0 is the initial position of the kink's center, and the breather, $\phi_{br}(x,t) = 4 \tan^{-1} \{ \sqrt{(1-\omega^2)/\omega^2} \sin[\omega\gamma\{t-v(x-x_0)\}] \operatorname{sech}[\gamma\sqrt{1-\omega^2}(x-x_0-vt)] \}$, where ω is the frequency of its internal oscillations ($0 < \omega < 1$).

The results of the interaction of the kink with the curvature in model A are shown in Fig. 1 [17]. The curvature variable f , initialized with a small random value, performs smooth oscillations with a frequency of $\omega \approx 2.866$. Notice that this is natural in this case, since the kink has an approximately fixed "mass," $M_k = \int_{-\infty}^{+\infty} u_x^2 dx$, which can be found to be $M_k = 8.247$ for the velocity $v \approx 0.25$. Then, Eq. (7) pre-

dicts the frequency of these oscillations $\omega \approx \sqrt{M_k} = 2.871$, which is very close to the numerically exact value. Fast small-amplitude oscillations of the kink's velocity, observed in Fig. 1 both with and without the curvature, are due to "hopping" over sites of a lattice (with spacing $h=0.1$) employed in the numerical scheme which solves Eq. (7). Notice, however, that in the top panel the mean velocity is ≈ 0.2503 , while in the bottom panel it is ≈ 0.2499 , hence the curvature oscillations increase the kink's velocity. This may be anticipated due to the presence of the positive definite factor $(1+f^2)$ in front of ϕ_{xx} in Eq. (6), which is expected to renormalize v^2 .

The curvature-breather interaction in model A is shown in Fig. 2. The frequency of the breather does not change significantly (it fluctuates between 0.87 and 0.93), but its amplitude decreases substantially (by more than 50%), resulting in its becoming much more mobile (the velocity increases to ≈ 0.425 from the initial value 0.25).

In model B, we examine the collision of a kink with a localized pulse of the substrate field f (for the breather, we have obtained results which are qualitatively similar to those presented below for the kink). We create the curvature pulse to the left or to the right of the kink. As ϕ_x vanishes far from the kink, Eq. (10) for f becomes a linear wave equation. Hence, we observe splitting of the pulse into left- and right-traveling ones. In the case where the kink is initially to the left of the pulse, it collides with the left-propagating fragment of the (split) pulse. In the opposite case, the kink is eventually caught by the copropagating right fragment of the (split) pulse. Numerical simulations shown in Fig. 3 demonstrate that the collision with the counterpropagating pulse

reduces the kink's velocity, while the interaction with the copropagating pulse gives rise to an increase of the velocity. In particular, in the former case, the mean speed of the kink after the collision is ≈ 0.2476 , while in the latter one, it increases to ≈ 0.2527 . Notice that in both cases a small fragment of the pulse that collides with the kink passes through it, while a larger fraction of the pulse is reflected by it.

In conclusion, we have proposed a field theory in the spirit of the Rahman-Parrinello molecular dynamics technique. The resulting equations couple the spatiotemporal evolution of the field to that of the underlying curvature of the space which carries the field. Coupled equations for the temporal or spatiotemporal evolution of the metric are obtained in a general setting, and, as an example, are solved together with the sine-Gordon field equation. The temporal evolution of the metric increases the velocity of the field solitons. The spatiotemporal evolution of the metric can increase or decrease their speed.

It would be particularly interesting to extend the models A and B to higher dimensions. It is also worth studying how the local evolution of the curvature affects kinematics and dynamics of the solitons, and to correlate such observations with the behavior of reactant chemical concentrations in chemical or biological environments with nontrivial geometry [18].

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